



Testing contextuality on quantum ensembles with one clean qubit; or: how I learned to stop worrying and love (use) my quantum computer.

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1 Contextuality

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- primer
- testing contextuality on ensembles of molecular nuclear spins in the solid state

- Contextuality

Motivation(s)

The main question, in some sense, is

given some preparation, ρ , and some observables

A B C

where

- \blacksquare A and B are compatible;
- \blacksquare A and C are compatible; and
- \blacksquare B and C are not (necessarily) compatible,

does the outcome of observable $A,\,\nu(A),$ depend on whether one measures B or C simultaneously with A?

- Contextuality
 - Motivation(s)

Other motivations

- Quantum computers seem to be able to solve some problems more efficiently than their classical counterparts — how come?
- What are small quantum computers good for?

- Contextuality
 - DIY (a simple proof)

DIY: contextuality

• consider a grid of 3x3 squares that can take on the values ± 1



• Evaluate the products of the rows and columns, and let $\alpha = \pi_{r_1} + \pi_{r_2} + \pi_{r_3} + \pi_{c_1} + \pi_{c_2} - \pi_{c_3}$

- Contextuality

DIY (a simple proof)

DIY: contextuality

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Figure: A histogram of α for all possible 2^9 configurations.

- Contextuality

└─ DIY (a simple proof)

DIY: contextuality - in other words

• for $\alpha = 6$, constrain the products to be:

	c_1	c_2	c_3	Π
r_1	?	?	?	+1
r_2	?	?	?	+1
r_3	?	?	?	+1

П	+1	+1	$^{-1}$
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are there any satisfying assignments?

- Contextuality

BKSC inequality

$\mathsf{Bell}^1 \ \mathsf{-Kochen}{-}\mathsf{Specker}^2 \ \mathsf{-}\mathsf{Cabello}^3 \ (\mathsf{BKSC}) \ inequality$

(for the Peres⁴ -Mermin⁵ square)

- Consider a system of 2 qubits prepared according to ρ.
- Measure the correlations between the observables listed in the columns and rows of the following table:

	c_1	c_2	c_3	
r_1	$\mathbb{Z}1$	$1\mathbb{Z}$	$\mathbb{Z}\mathbb{Z}$	
r_2	1X	X1	XX	
r_3	ZX	XZ	YY	



	П	+1	+1	-1
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- Evaluate $\beta = \langle r_1 \rangle + \langle r_2 \rangle + \langle r_3 \rangle + \langle c_1 \rangle + \langle c_2 \rangle \langle c_3 \rangle$
- where, e.g., $\langle r_1 \rangle = \langle \pi_{\{\mathbb{Z}\,\mathbb{1},\mathbb{1}\mathbb{Z},\mathbb{Z}\mathbb{Z}\}} \rangle = \langle \mathbb{Z}\mathbb{1} \cdot \mathbb{1}\mathbb{Z} \cdot \mathbb{Z}\mathbb{Z} \rangle_{\rho}$, and so forth.
- then, for any state ρ :
 - **•** For any NCHV theory: $\beta \leq 4$
 - **Quantum Mechanics:** $\beta = 6$

¹J. S. Bell. Rev. Mod. Phys. 38, 447-452 (1966).

- ²S. Kochen, and E. P. Specker. J. Math. Mech. 17, 59-87 (1967).
- ³A. Cabello, *Phys. Rev. Lett.* **101**, 210401 (2008).
- ⁴A. Peres. *Physics Letters A*. **151**, 107 (1990).
- ⁵N. D. Mermin. Phys. Rev. Lett. 65, 3373 (1990).

- Contextuality
 - Experimental efforts

Experimental efforts

Two trapped Ca ions ⁶



- $\begin{array}{l} \mathbf{b} & U[\sigma_{s} \oplus \sigma_{s}] = U_{Y}^{(S)}(\neg \pi/2) \; U_{s}^{(1)}(\pi/2) & U[\sigma_{s} \oplus \sigma_{s}] = U_{X}^{(S)}(\neg \pi/2) \; U_{s}(\pi/2) \\ & U[\sigma_{s} \oplus \sigma_{s}] = U_{X}^{(S)}(\neg \pi/2) \; U_{s}^{(1)}(\pi/2) & U[\sigma_{s} \oplus \sigma_{s}] = U_{Y}^{(S)}(\neg \pi/2) \; U_{s}(\neg \pi/2) \\ & U[\sigma_{s} \oplus \sigma_{s}] = U_{Y}^{(S)}(\neg \pi/2) \; U_{s}^{(1)}(\pi/2) \; U_{Y}(\pi/2) \\ \end{array}$
- Path and polarization degrees of freedom of a single photon ⁷



⁶G. Kirchmair et al., Nature **460**, 494 (2009).

⁷E. Amselem et al., Phys. Rev. Lett. 103, 160405 (2009).

A bit of quantum computer science

Alternatively,

- Introduce a probe two-level system (prepare in +1 eigenstate of \mathbb{X})
- To measure any two-outcome observable S:
 - if the system is in a +1 eigenstate of S, do nothing; and
 - if it is in a -1 eigenstate, apply a phase flip (pauli- \mathbb{Z}) to the probe qubit.
- For $S = P_+ P_-$, where P_+ and P_- are the projectors on the +1 and -1 eigenspaces of S, define

$$U_S = \mathbb{1}_2 \otimes P_+ + \mathbb{Z} \otimes P_-$$

This transformation can also be expressed as a controlled operation dependent on the state of the probe qubit:

$$U_S = \mathbb{1}_2 \otimes P_+ + \mathbb{Z} \otimes P_-$$

= $\frac{1}{2} (\mathbb{1}_2 + \mathbb{Z}) \otimes (P_+ + P_-) + \frac{1}{2} (\mathbb{1}_2 - \mathbb{Z}) \otimes (P_+ - P_-)$
= $|0\rangle \langle 0| \otimes \mathbb{1}_d + |1\rangle \langle 1| \otimes S$,

which is unitary for S unitary.

- A bit of quantum computer science
 - circuit for ensemble measurement of correlations

Ensemble measurement of $\langle S_1 S_2 S_3 \rangle_{1}$



* An ensemble measurement of $\ensuremath{\mathbb{X}}$ on the probe qubit evaluates to the required correlation.

* Incidentally, this model of computation is known as Deterministic Quantum Computation with one qubit (DQC1)

X

X

X)

 \mathbb{Z}

 \mathbb{X}

X

 \mathbb{Z}

Y

Y

A bit of quantum computer science

└─ circuit for ensemble measurement of correlations



Solid-state NMR for QIP



Figure: Malonic acid $(C_3H_4O_4)$ molecule and Hamiltonian parameters (all values in kHz). Elements along the diagonal represent chemical shifts, ω_i , with respect to the transmitter frequency (with the Hamiltonian $\sum_i \pi \omega_i \mathbb{Z}_i$). Above the diagonal are dipolar coupling constants ($\sum_{i < j} \pi D_{i,j} (2 \mathbb{Z}_i \mathbb{Z}_j - \mathbb{X}_i \mathbb{X}_j - \mathbb{Y}_i \mathbb{Y}_j)$, and below the diagonal are J coupling constants, ($\sum_{i < j} \frac{\pi}{2} J_{i,j} (\mathbb{Z}_i \mathbb{Z}_j + \mathbb{X}_i \mathbb{X}_j + \mathbb{Y}_i \mathbb{Y}_j)$). An accurate natural Hamiltonian is necessary for high fidelity control and is obtained from precise spectral fitting.

Solid-state NMR for QIP

└─ testing contextuality on ensembles of molecular nuclear spins in the solid state

Experimental results: $\beta = 5.2 \pm 0.1$



Figure: Summary of experimental results (solid lines) and the corresponding spectral fits (dashed lines). Shown are (a) (in green) a proton-decoupled ^{13}C spectrum following polarization-transfer from the abundant protons; (b) (in blue) the spectrum produced by the initial preparation procedure, $\mathbb{X} \otimes \frac{\mathbb{1}_2}{2} \otimes \frac{\mathbb{1}_2}{2}$; and (c) (in red) the average of the six spectra corresponding to the six terms in β weighted by the appropriate signs.

Solid-state NMR for QIP

testing contextuality on ensembles of molecular nuclear spins in the solid state

What I was trying to say is

- BKSC inequality creates a separation between NCHV theories and QM.
 - Non-contextuality is testable ...
 - using DQC1
 - ... on ensembles
 - ... of nuclear spins
 - ... in the solid state
 - using their magnetic resonance.

Quantum Information Processing is useful for tasks other than factoring.

Solid-state NMR for QIP

testing contextuality on ensembles of molecular nuclear spins in the solid state

If you have been,

Thanks for listening.