# Testing contextuality on quantum ensembles with one clean qubit; or: how I learned to stop worrying and love (use) my quantum computer. 

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1 Contextuality

- Motivation(s)
- DIY (a simple proof)
- BKSC inequality
- Experimental efforts

2 A bit of quantum computer science
■ circuit for ensemble measurement of correlations

3 Solid-state NMR for QIP

- primer
- testing contextuality on ensembles of molecular nuclear spins in the solid state

The main question, in some sense, is
given some preparation, $\rho$, and some observables
$A \quad B$
where

- $A$ and $B$ are compatible;
- $A$ and $C$ are compatible; and
- $B$ and $C$ are not (necessarily) compatible,
does the outcome of observable $A, \nu(A)$, depend on whether one measures $B$ or $C$ simultaneously with $A$ ?


## Other motivations

- Quantum computers seem to be able to solve some problems more efficiently than their classical counterparts - how come?
■ What are small quantum computers good for?


## DIY: contextuality

- consider a grid of $3 \times 3$ squares that can take on the values $\pm 1$

- Evaluate the products of the rows and columns, and let $\alpha=\pi_{r_{1}}+\pi_{r_{2}}+\pi_{r_{3}}+\pi_{c_{1}}+\pi_{c_{2}}-\pi_{c_{3}}$


## DIY: contextuality

- consider a grid of $3 \times 3$ squares that can take on the values $\pm 1$

|  | $c_{1}$ | $c_{2}$ | $c_{3}$ | $\prod$ |
| :---: | :---: | :---: | :---: | :---: |
| $r_{1}$ | $\pm 1$ | $\pm 1$ | $\pm 1$ | $\pm 1$ |
| $r_{2}$ | $\pm 1$ | $\pm 1$ | $\pm 1$ | $\pm 1$ |
| $r_{3}$ | $\pm 1$ | $\pm 1$ | $\pm 1$ | $\pm 1$ |
| $\Pi$ | $\pm 1$ | $\pm 1$ | $\pm 1$ |  |

- Evaluate the products of the rows and columns, and let $\alpha=\pi_{r_{1}}+\pi_{r_{2}}+\pi_{r_{3}}+\pi_{c_{1}}+\pi_{c_{2}}-\pi_{c_{3}}$


Figure: A histogram of $\alpha$ for all possible $2^{9}$ configurations.

## DIY: contextuality - in other words

- for $\alpha=6$, constrain the products to be:


■ are there any satisfying assignments?

## Bell ${ }^{1}$-Kochen-Specker ${ }^{2}$ - Cabello ${ }^{3}$ (BKSC) inequality

 (for the Peres ${ }^{4}$-Mermin ${ }^{5}$ square)- Consider a system of 2 qubits prepared according to $\rho$.
■ Measure the correlations between the observables listed in the columns and rows of the following table:

|  | $c_{1}$ | $c_{2}$ | $c_{3}$ |
| :---: | :---: | :---: | :---: |
| $r_{1}$ | $\mathbb{Z} \mathbb{1}$ | $\mathbb{1} \mathbb{Z}$ | $\mathbb{Z} \mathbb{Z}$ |
| $r_{2}$ | $\mathbb{1} \mathbb{X}$ | $\mathbb{X} \mathbb{1}$ | $\mathbb{X} \mathbb{X}$ |
| $r_{3}$ | $\mathbb{Z} \mathbb{X}$ | $\mathbb{X} \mathbb{Z}$ | $\mathbb{Y} \mathbb{Y}$ |
| $+\mathbb{1}$ |  |  |  |
| $+\mathbb{1}$ |  |  |  |
| $+\mathbb{1}$ |  |  |  |


| $\Pi$ | $+\mathbb{1}$ | $+\mathbb{1}$ | $-\mathbb{1}$ |
| :--- | :--- | :--- | :--- |

■ Evaluate $\beta=\left\langle r_{1}\right\rangle+\left\langle r_{2}\right\rangle+\left\langle r_{3}\right\rangle+\left\langle c_{1}\right\rangle+\left\langle c_{2}\right\rangle-\left\langle c_{3}\right\rangle$
$■$ where, e.g., $\left\langle r_{1}\right\rangle=\left\langle\pi_{\{\mathbb{Z}, \mathbb{Z}, \mathbb{Z} \mathbb{Z}\}}\right\rangle=\langle\mathbb{Z} \mathbb{1} \cdot \mathbb{1} \mathbb{Z} \cdot \mathbb{Z} \mathbb{Z}\rangle_{\rho}$, and so forth.

- then, for any state $\rho$ :
- For any NCHV theory: $\beta \leq 4$
- Quantum Mechanics: $\beta=6$

[^0]
## Experimental efforts

- Two trapped Ca ions ${ }^{6}$
a

b

$$
\begin{aligned}
& U\left[\sigma_{y} 8 \sigma_{y}\right]=U_{X}^{M S}(-\pi / 2) U_{z}^{(1)}(\pi / 2) \quad U\left[\sigma_{z} \otimes \sigma_{x}\right]=U_{Y}^{M S}(-\pi / 2) U_{X}(-\pi / 2) \\
& U\left[\sigma_{z} \sigma_{]}=U_{y}^{1 s}(-\pi / 2) U_{z}^{11(\pi / 2)} U_{\chi \pi / 2}\right)
\end{aligned}
$$

- Path and polarization degrees of freedom of a single photon ${ }^{7}$


[^1]Alternatively,
■ Introduce a probe two-level system (prepare in +1 eigenstate of $\mathbb{X}$ )
■ To measure any two-outcome observable $S$ :

- if the system is in a +1 eigenstate of $S$, do nothing; and
- if it is in a -1 eigenstate, apply a phase flip (pauli- $\mathbb{Z}$ ) to the probe qubit.

■ For $S=P_{+}-P_{-}$, where $P_{+}$and $P_{-}$are the projectors on the +1 and -1 eigenspaces of $S$, define

$$
U_{S}=\mathbb{1}_{2} \otimes P_{+}+\mathbb{Z} \otimes P_{-}
$$

This transformation can also be expressed as a controlled operation dependent on the state of the probe qubit:

$$
\begin{aligned}
U_{S} & =\mathbb{1}_{2} \otimes P_{+}+\mathbb{Z} \otimes P_{-} \\
& =\frac{1}{2}\left(\mathbb{1}_{2}+\mathbb{Z}\right) \otimes\left(P_{+}+P_{-}\right)+\frac{1}{2}\left(\mathbb{1}_{2}-\mathbb{Z}\right) \otimes\left(P_{+}-P_{-}\right) \\
& =|0\rangle\langle 0| \otimes \mathbb{1}_{d}+|1\rangle\langle 1| \otimes S,
\end{aligned}
$$

which is unitary for $S$ unitary.

## Ensemble measurement of $\left\langle S_{1} S_{2} S_{3}\right\rangle_{\mathbb{1}}$



* An ensemble measurement of $\mathbb{X}$ on the probe qubit evaluates to the required correlation.
* Incidentally, this model of computation is known as Deterministic Quantum

Computation with one qubit (DQC1)

$\left(C_{1}\right)$

$\left(r_{2}\right)$
$\left(C_{2}\right)$

$\left(r_{3}\right)$

$\left(C_{3}\right)$



| kHz | $\mathrm{C}_{1}$ | $\mathrm{C}_{2}$ | $\mathrm{C}_{m}$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{C}_{1}$ | 6.380 | 0.297 | 0.780 |
| $\mathrm{C}_{2}$ | -0.025 | -1.533 | 1.050 |
| $\mathrm{C}_{m}$ | 0.071 | 0.042 | -5.650 |



Figure: Malonic acid $\left(\mathrm{C}_{3} \mathrm{H}_{4} \mathrm{O}_{4}\right)$ molecule and Hamiltonian parameters (all values in kHz ). Elements along the diagonal represent chemical shifts, $\omega_{i}$, with respect to the transmitter frequency (with the Hamiltonian $\sum_{i} \pi \omega_{i} \mathbb{Z}_{i}$ ). Above the diagonal are dipolar coupling constants $\left(\sum_{i<j} \pi D_{i, j}\left(2 \mathbb{Z}_{i} \mathbb{Z}_{j}-\mathbb{X}_{i} \mathbb{X}_{j}-\mathbb{Y}_{i} \mathbb{Y}_{j}\right)\right.$, and below the diagonal are J coupling constants, $\left(\sum_{i<j} \frac{\pi}{2} J_{i, j}\left(\mathbb{Z}_{i} \mathbb{Z}_{j}+\mathbb{X}_{i} \mathbb{X}_{j}+\mathbb{Y}_{i} \mathbb{Y}_{j}\right)\right.$. An accurate natural Hamiltonian is necessary for high fidelity control and is obtained from precise spectral fitting.

## Experimental results: $\beta=5.2 \pm 0.1$



Figure: Summary of experimental results (solid lines) and the corresponding spectral fits (dashed lines). Shown are (a) (in green) a proton-decoupled ${ }^{13} \mathrm{C}$ spectrum following polarization-transfer from the abundant protons; (b) (in blue) the spectrum produced by the initial preparation procedure, $\mathbb{X} \otimes \frac{\mathbb{1}_{2}}{2} \otimes \frac{\mathbb{1}_{2}}{2}$; and (c) (in red) the average of the six spectra corresponding to the six terms in $\beta$ weighted by the appropriate signs.

## What I was trying to say is

■ BKSC inequality creates a separation between NCHV theories and QM.
■ Non-contextuality is testable ...

- ... using DQC1
- ... on ensembles
- ... of nuclear spins
- ... in the solid state
- ... using their magnetic resonance.

■ Quantum Information Processing is useful for tasks other than factoring.

## If you have been,

## Thanks for listening.


[^0]:    ${ }^{1}$ J. S. Bell.Rev. Mod. Phys. 38, 447-452 (1966).
    ${ }^{2}$ S. Kochen, and E. P. Specker. J. Math. Mech. 17, 59-87 (1967).
    ${ }^{3}$ A. Cabello, Phys. Rev. Lett. 101, 210401 (2008).
    ${ }^{4}$ A. Peres. Physics Letters A. 151, 107 (1990).
    ${ }^{5}$ N. D. Mermin. Phys. Rev. Lett. 65, 3373 (1990).

[^1]:    ${ }^{6}$ G. Kirchmair et al., Nature 460, 494 (2009).
    ${ }^{7}$ E. Amselem et al., Phys. Rev. Lett. 103, 160405 (2009).

